

Top-quark mass measurement using events with missing transverse energy and jets at CDF

T. Aaltonen,²¹ B. Álvarez González^{w,9} S. Amerio,⁴¹ D. Amidei,³² A. Anastassov,³⁶ A. Annovi,¹⁷ J. Antos,¹² G. Apollinari,¹⁵ J.A. Appel,¹⁵ A. Apresyan,⁴⁶ T. Arisawa,⁵⁶ A. Artikov,¹³ J. Asaadi,⁵¹ W. Ashmanskas,¹⁵ B. Auerbach,⁵⁹ A. Aurisano,⁵¹ F. Azfar,⁴⁰ W. Badgett,¹⁵ A. Barbaro-Galtieri,²⁶ V.E. Barnes,⁴⁶ B.A. Barnett,²³ P. Barria^{dd,44} P. Bartos,¹² M. Bauce^{bb,41} G. Bauer,³⁰ F. Bedeschi,⁴⁴ D. Beecher,²⁸ S. Behari,²³ G. Bellettini^{cc,44} J. Bellinger,⁵⁸ D. Benjamin,¹⁴ A. Beretvas,¹⁵ A. Bhatti,⁴⁸ M. Binkley^{*,15} D. Bisello^{bb,41} I. Bizjak^{hh,28} K.R. Bland,⁵ B. Blumenfeld,²³ A. Bocci,¹⁴ A. Bodek,⁴⁷ D. Bortoletto,⁴⁶ J. Boudreau,⁴⁵ A. Boveia,¹¹ L. Brigliadori^{aa,6} A. Brisuda,¹² C. Bromberg,³³ E. Brucken,²¹ M. Bucchiantonio^{cc,44} J. Budagov,¹³ H.S. Budd,⁴⁷ S. Budd,²² K. Burkett,¹⁵ G. Busetto^{bb,41} P. Bussey,¹⁹ A. Buzatu,³¹ C. Calancha,²⁹ S. Camarda,⁴ M. Campanelli,²⁸ M. Campbell,³² F. Canelli^{11,15} B. Carls,²² D. Carlsmith,⁵⁸ R. Carosi,⁴⁴ S. Carrillo^{k,16} S. Carron,¹⁵ B. Casal,⁹ M. Casarsa,¹⁵ A. Castro^{aa,6} P. Catastini,²⁰ D. Cauz,⁵² V. Cavaliere,²² M. Cavalli-Sforza,⁴ A. Cerri^{e,26} L. Cerrito^{q,28} Y.C. Chen,¹ M. Chertok,⁷ G. Chiarelli,⁴⁴ G. Chlachidze,¹⁵ F. Chlebana,¹⁵ K. Cho,²⁵ D. Chokheli,¹³ J.P. Chou,²⁰ W.H. Chung,⁵⁸ Y.S. Chung,⁴⁷ C.I. Ciobanu,⁴² M.A. Ciocci^{dd,44} A. Clark,¹⁸ C. Clarke,⁵⁷ G. Compostella^{bb,41} M.E. Convery,¹⁵ J. Conway,⁷ M. Corbo,⁴² M. Cordelli,¹⁷ C.A. Cox,⁷ D.J. Cox,⁷ F. Crescioli^{cc,44} C. Cuenca Almenar,⁵⁹ J. Cuevas^{w,9} R. Culbertson,¹⁵ D. Dagenhart,¹⁵ N. d'Ascenzo^{u,42} M. Datta,¹⁵ P. de Barbaro,⁴⁷ S. De Cecco,⁴⁹ G. De Lorenzo,⁴ M. Dell'Orso^{cc,44} C. Deluca,⁴ L. Demortier,⁴⁸ J. Deng^{b,14} M. Deninno,⁶ F. Devoto,²¹ M. d'Errico^{bb,41} A. Di Canto^{cc,44} B. Di Ruzza,⁴⁴ J.R. Dittmann,⁵ M. D'Onofrio,²⁷ S. Donati^{cc,44} P. Dong,¹⁵ M. Dorigo,⁵² T. Dorigo,⁴¹ K. Ebina,⁵⁶ A. Elagin,⁵¹ A. Eppig,³² R. Erbacher,⁷ D. Errede,²² S. Errede,²² N. Ershaidat^{z,42} R. Eusebi,⁵¹ H.C. Fang,²⁶ S. Farrington,⁴⁰ M. Feindt,²⁴ J.P. Fernandez,²⁹ C. Ferrazza^{ee,44} R. Field,¹⁶ G. Flanagan^{s,46} R. Forrest,⁷ M.J. Frank,⁵ M. Franklin,²⁰ J.C. Freeman,¹⁵ Y. Funakoshi,⁵⁶ I. Furic,¹⁶ M. Gallinaro,⁴⁸ J. Galyardt,¹⁰ J.E. Garcia,¹⁸ A.F. Garfinkel,⁴⁶ P. Garosi^{dd,44} H. Gerberich,²² E. Gerchtein,¹⁵ S. Giagu^{ff,49} V. Giakoumopoulou,³ P. Giannetti,⁴⁴ K. Gibson,⁴⁵ C.M. Ginsburg,¹⁵ N. Giokaris,³ P. Giromini,¹⁷ M. Giunta,⁴⁴ G. Giurgiu,²³ V. Glagolev,¹³ D. Glenzinski,¹⁵ M. Gold,³⁵ D. Goldin,⁵¹ N. Goldschmidt,¹⁶ A. Golossanov,¹⁵ G. Gomez,⁹ G. Gomez-Ceballos,³⁰ M. Goncharov,³⁰ O. González,²⁹ I. Gorelov,³⁵ A.T. Goshaw,¹⁴ K. Goulianos,⁴⁸ S. Grinstein,⁴ C. Grosso-Pilcher,¹¹ R.C. Group^{55,15} J. Guimaraes da Costa,²⁰ Z. Gunay-Unalan,³³ C. Haber,²⁶ S.R. Hahn,¹⁵ E. Halkiadakis,⁵⁰ A. Hamaguchi,³⁹ J.Y. Han,⁴⁷ F. Happacher,¹⁷ K. Hara,⁵³ D. Hare,⁵⁰ M. Hare,⁵⁴ R.F. Harr,⁵⁷ K. Hatakeyama,⁵ C. Hays,⁴⁰ M. Heck,²⁴ J. Heinrich,⁴³ M. Herndon,⁵⁸ S. Hewamanage,⁵ D. Hidas,⁵⁰ A. Hocker,¹⁵ W. Hopkins^{f,15} D. Horn,²⁴ S. Hou,¹ R.E. Hughes,³⁷ M. Hurwitz,¹¹ U. Husemann,⁵⁹ N. Hussain,³¹ M. Hussein,³³ J. Huston,³³ G. Introzzi,⁴⁴ M. Iori^{ff,49} A. Ivanov^{o,7} E. James,¹⁵ D. Jang,¹⁰ B. Jayatilaka,¹⁴ E.J. Jeon,²⁵ M.K. Jha,⁶ S. Jindariani,¹⁵ W. Johnson,⁷ M. Jones,⁴⁶ K.K. Joo,²⁵ S.Y. Jun,¹⁰ T.R. Junk,¹⁵ T. Kamon,⁵¹ P.E. Karchin,⁵⁷ A. Kasmi,⁵ Y. Kato^{n,39} W. Ketchum,¹¹ J. Keung,⁴³ V. Khotilovich,⁵¹ B. Kilminster,¹⁵ D.H. Kim,²⁵ H.S. Kim,²⁵ H.W. Kim,²⁵ J.E. Kim,²⁵ M.J. Kim,¹⁷ S.B. Kim,²⁵ S.H. Kim,⁵³ Y.K. Kim,¹¹ N. Kimura,⁵⁶ M. Kirby,¹⁵ S. Klimenko,¹⁶ K. Kondo^{*,56} D.J. Kong,²⁵ J. Konigsberg,¹⁶ A.V. Kotwal,¹⁴ M. Kreps,²⁴ J. Kroll,⁴³ D. Krop,¹¹ N. Krumnack^{l,5} M. Kruse,¹⁴ V. Krutelyov^{c,51} T. Kuhr,²⁴ M. Kurata,⁵³ S. Kwang,¹¹ A.T. Laasanen,⁴⁶ S. Lami,⁴⁴ S. Lammel,¹⁵ M. Lancaster,²⁸ R.L. Lander,⁷ K. Lannon^{v,37} A. Lath,⁵⁰ G. Latino^{cc,44} T. LeCompte,² E. Lee,⁵¹ H.S. Lee,¹¹ J.S. Lee,²⁵ S.W. Lee^{x,51} S. Leo^{cc,44} S. Leone,⁴⁴ J.D. Lewis,¹⁵ A. Limosani^{r,14} C.-J. Lin,²⁶ J. Linacre,⁴⁰ M. Lindgren,¹⁵ E. Lipeles,⁴³ A. Lister,¹⁸ D.O. Litvintsev,¹⁵ C. Liu,⁴⁵ Q. Liu,⁴⁶ T. Liu,¹⁵ S. Lockwitz,⁵⁹ A. Loginov,⁵⁹ D. Lucchesi^{bb,41} J. Lueck,²⁴ P. Lujan,²⁶ P. Lukens,¹⁵ G. Lungu,⁴⁸ J. Lys,²⁶ R. Lysak,¹² R. Madrak,¹⁵ K. Maeshima,¹⁵ K. Makhoul,³⁰ S. Malik,⁴⁸ G. Manca^{a,27} A. Manousakis-Katsikakis,³ F. Margaroli,⁴⁶ C. Marino,²⁴ M. Martínez,⁴ R. Martínez-Ballarín,²⁹ P. Mastrandrea,⁴⁹ M.E. Mattson,⁵⁷ P. Mazzanti,⁶ K.S. McFarland,⁴⁷ P. McIntyre,⁵¹ R. McNulty^{i,27} A. Mehta,²⁷ P. Mehtala,²¹ A. Menzione,⁴⁴ C. Mesropian,⁴⁸ T. Miao,¹⁵ D. Mietlicki,³² A. Mitra,¹ H. Miyake,⁵³ S. Moed,²⁰ N. Moggi,⁶ M.N. Mondragon^{k,15} C.S. Moon,²⁵ R. Moore,¹⁵ M.J. Morello,¹⁵ J. Morlock,²⁴ P. Movilla Fernandez,¹⁵ A. Mukherjee,¹⁵ Th. Muller,²⁴ P. Murat,¹⁵ M. Mussini^{aa,6} J. Nachtman^{m,15} Y. Nagai,⁵³ J. Naganoma,⁵⁶ I. Nakano,³⁸ A. Napier,⁵⁴ J. Nett,⁵¹ C. Neu,⁵⁵ M.S. Neubauer,²² J. Nielsen^{d,26} L. Nodulman,² O. Norniella,²² E. Nurse,²⁸ L. Oakes,⁴⁰ S.H. Oh,¹⁴ Y.D. Oh,²⁵ I. Oksuzian,⁵⁵ T. Okusawa,³⁹ R. Orava,²¹ L. Ortolan,⁴ S. Pagan Griso^{bb,41} C. Pagliarone,⁵² E. Palencia^{e,9} V. Papadimitriou,¹⁵ A.A. Paramonov,² J. Patrick,¹⁵ G. Pauletta^{gg,52} M. Paulini,¹⁰ C. Paus,³⁰ D.E. Pellett,⁷ A. Penzo,⁵² T.J. Phillips,¹⁴ G. Piacentino,⁴⁴ E. Pianori,⁴³ J. Pilot,³⁷ K. Pitts,²² C. Plager,⁸ L. Pondrom,⁵⁸ K. Potamianos,⁴⁶ O. Poukhov^{*,13} F. Prokoshin^{y,13} A. Pronko,¹⁵ F. Ptohos^{q,17} E. Pueschel,¹⁰ G. Punzi^{cc,44} J. Pursley,⁵⁸ A. Rahaman,⁴⁵ V. Ramakrishnan,⁵⁸ N. Ranjan,⁴⁶ I. Redondo,²⁹ P. Renton,⁴⁰ M. Rescigno,⁴⁹ T. Riddick,²⁸ F. Rimondi^{aa,6} L. Ristori^{44,15} A. Robson,¹⁹

T. Rodrigo,⁹ T. Rodriguez,⁴³ E. Rogers,²² S. Rolli^h,⁵⁴ R. Roser,¹⁵ M. Rossi,⁵² F. Rubbo,¹⁵ F. Ruffini^{dd},⁴⁴
A. Ruiz,⁹ J. Russ,¹⁰ V. Rusu,¹⁵ A. Safonov,⁵¹ W.K. Sakumoto,⁴⁷ Y. Sakurai,⁵⁶ L. Santi^{gg},⁵² L. Sartori,⁴⁴
K. Sato,⁵³ V. Saveliev^u,⁴² A. Savoy-Navarro,⁴² P. Schlabach,¹⁵ A. Schmidt,²⁴ E.E. Schmidt,¹⁵ M.P. Schmidt^{*},⁵⁹
M. Schmitt,³⁶ T. Schwarz,⁷ L. Scodellaro,⁹ A. Scribano^{dd},⁴⁴ F. Scuri,⁴⁴ A. Sedov,⁴⁶ S. Seidel,³⁵ Y. Seiya,³⁹
A. Semenov,¹³ F. Sforza^{cc},⁴⁴ A. Sfyrila,²² S.Z. Shalhout,⁷ T. Shears,²⁷ P.F. Shepard,⁴⁵ M. Shimojima^t,⁵³
S. Shiraishi,¹¹ M. Shochet,¹¹ I. Shreyber,³⁴ A. Simonenko,¹³ P. Sinervo,³¹ A. Sissakian^{*},¹³ K. Sliwa,⁵⁴ J.R. Smith,⁷
F.D. Snider,¹⁵ A. Soha,¹⁵ S. Somalwar,⁵⁰ V. Sorin,⁴ P. Squillacioti,⁴⁴ M. Stancari,¹⁵ M. Stanitzki,⁵⁹
R. St. Denis,¹⁹ B. Stelzer,³¹ O. Stelzer-Chilton,³¹ D. Stentz,³⁶ J. Strologas,³⁵ G.L. Strycker,³² Y. Sudo,⁵³
A. Sukhanov,¹⁶ I. Suslov,¹³ K. Takemasa,⁵³ Y. Takeuchi,⁵³ J. Tang,¹¹ M. Tecchio,³² P.K. Teng,¹ J. Thom^f,¹⁵
J. Thome,¹⁰ G.A. Thompson,²² E. Thomson,⁴³ P. Ttito-Guzmán,²⁹ S. Tkaczyk,¹⁵ D. Toback,⁵¹ S. Tokar,¹²
K. Tollefson,³³ T. Tomura,⁵³ D. Tonelli,¹⁵ S. Torre,¹⁷ D. Torretta,¹⁵ P. Totaro,⁴¹ M. Trovato^{ee},⁴⁴ Y. Tu,⁴³
F. Ukegawa,⁵³ S. Uozumi,²⁵ A. Varganov,³² F. Vázquez^k,¹⁶ G. Velez,¹⁵ C. Vellidis,³ M. Vidal,²⁹ I. Vila,⁹
R. Vilar,⁹ J. Vizán,⁹ M. Vogel,³⁵ G. Volpi^{cc},⁴⁴ P. Wagner,⁴³ R.L. Wagner,¹⁵ T. Wakisaka,³⁹ R. Wallny,⁸
S.M. Wang,¹ A. Warburton,³¹ D. Waters,²⁸ M. Weinberger,⁵¹ W.C. Wester III,¹⁵ B. Whitehouse,⁵⁴ D. Whiteson^b,⁴³
A.B. Wicklund,² E. Wicklund,¹⁵ S. Wilbur,¹¹ F. Wick,²⁴ H.H. Williams,⁴³ J.S. Wilson,³⁷ P. Wilson,¹⁵ B.L. Winer,³⁷
P. Wittich^g,¹⁵ S. Wolbers,¹⁵ H. Wolfe,³⁷ T. Wright,³² X. Wu,¹⁸ Z. Wu,⁵ K. Yamamoto,³⁹ J. Yamaoka,¹⁴
T. Yang,¹⁵ U.K. Yang^p,¹¹ Y.C. Yang,²⁵ W.-M. Yao,²⁶ G.P. Yeh,¹⁵ K. Yi^m,¹⁵ J. Yoh,¹⁵ K. Yorita,⁵⁶
T. Yoshida^j,³⁹ G.B. Yu,¹⁴ I. Yu,²⁵ S.S. Yu,¹⁵ J.C. Yun,¹⁵ A. Zanetti,⁵² Y. Zeng,¹⁴ and S. Zucchelli^{aa6}

(CDF Collaboration[†])

¹*Institute of Physics, Academia Sinica, Taipei, Taiwan 11529, Republic of China*

²*Argonne National Laboratory, Argonne, Illinois 60439, USA*

³*University of Athens, 157 71 Athens, Greece*

⁴*Institut de Física d'Altes Energies, ICREA, Universitat Autònoma de Barcelona, E-08193, Bellaterra (Barcelona), Spain*

⁵*Baylor University, Waco, Texas 76798, USA*

⁶*Istituto Nazionale di Fisica Nucleare Bologna, ^{aa}University of Bologna, I-40127 Bologna, Italy*

⁷*University of California, Davis, Davis, California 95616, USA*

⁸*University of California, Los Angeles, Los Angeles, California 90024, USA*

⁹*Instituto de Física de Cantabria, CSIC-University of Cantabria, 39005 Santander, Spain*

¹⁰*Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA*

¹¹*Enrico Fermi Institute, University of Chicago, Chicago, Illinois 60637, USA*

¹²*Comenius University, 842 48 Bratislava, Slovakia; Institute of Experimental Physics, 040 01 Kosice, Slovakia*

¹³*Joint Institute for Nuclear Research, RU-141980 Dubna, Russia*

¹⁴*Duke University, Durham, North Carolina 27708, USA*

¹⁵*Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA*

¹⁶*University of Florida, Gainesville, Florida 32611, USA*

¹⁷*Laboratori Nazionali di Frascati, Istituto Nazionale di Fisica Nucleare, I-00044 Frascati, Italy*

¹⁸*University of Geneva, CH-1211 Geneva 4, Switzerland*

¹⁹*Glasgow University, Glasgow G12 8QQ, United Kingdom*

²⁰*Harvard University, Cambridge, Massachusetts 02138, USA*

²¹*Division of High Energy Physics, Department of Physics,
University of Helsinki and Helsinki Institute of Physics, FIN-00014, Helsinki, Finland*

²²*University of Illinois, Urbana, Illinois 61801, USA*

²³*The Johns Hopkins University, Baltimore, Maryland 21218, USA*

²⁴*Institut für Experimentelle Kernphysik, Karlsruhe Institute of Technology, D-76131 Karlsruhe, Germany*

²⁵*Center for High Energy Physics: Kyungpook National University,*

Daegu 702-701, Korea; Seoul National University, Seoul 151-742,

Korea; Sungkyunkwan University, Suwon 440-746,

Korea; Korea Institute of Science and Technology Information,

Daejeon 305-806, Korea; Chonnam National University, Gwangju 500-757,

Korea; Chonbuk National University, Jeonju 561-756, Korea

²⁶*Ernest Orlando Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

²⁷*University of Liverpool, Liverpool L69 7ZE, United Kingdom*

²⁸*University College London, London WC1E 6BT, United Kingdom*

²⁹*Centro de Investigaciones Energeticas Medioambientales y Tecnológicas, E-28040 Madrid, Spain*

³⁰*Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA*

³¹*Institute of Particle Physics: McGill University, Montréal, Québec,*

Canada H3A 2T8; Simon Fraser University, Burnaby, British Columbia,

Canada V5A 1S6; University of Toronto, Toronto, Ontario,

Canada M5S 1A7; and TRIUMF, Vancouver, British Columbia, Canada V6T 2A3

³²*University of Michigan, Ann Arbor, Michigan 48109, USA*

- ³³Michigan State University, East Lansing, Michigan 48824, USA
³⁴Institution for Theoretical and Experimental Physics, ITEP, Moscow 117259, Russia
³⁵University of New Mexico, Albuquerque, New Mexico 87131, USA
³⁶Northwestern University, Evanston, Illinois 60208, USA
³⁷The Ohio State University, Columbus, Ohio 43210, USA
³⁸Okayama University, Okayama 700-8530, Japan
³⁹Osaka City University, Osaka 588, Japan
⁴⁰University of Oxford, Oxford OX1 3RH, United Kingdom
⁴¹Istituto Nazionale di Fisica Nucleare, Sezione di Padova-Trento, ^{bb}University of Padova, I-35131 Padova, Italy
⁴²LPNHE, Université Pierre et Marie Curie/IN2P3-CNRS, UMR7585, Paris, F-75252 France
⁴³University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA
⁴⁴Istituto Nazionale di Fisica Nucleare Pisa, ^{cc}University of Pisa,
^{dd}University of Siena and ^{ee}Scuola Normale Superiore, I-56127 Pisa, Italy
⁴⁵University of Pittsburgh, Pittsburgh, Pennsylvania 15260, USA
⁴⁶Purdue University, West Lafayette, Indiana 47907, USA
⁴⁷University of Rochester, Rochester, New York 14627, USA
⁴⁸The Rockefeller University, New York, New York 10065, USA
⁴⁹Istituto Nazionale di Fisica Nucleare, Sezione di Roma 1,
^{ff}Sapienza Università di Roma, I-00185 Roma, Italy
⁵⁰Rutgers University, Piscataway, New Jersey 08855, USA
⁵¹Texas A&M University, College Station, Texas 77843, USA
⁵²Istituto Nazionale di Fisica Nucleare Trieste/Udine,
I-34100 Trieste, ^{gg}University of Udine, I-33100 Udine, Italy
⁵³University of Tsukuba, Tsukuba, Ibaraki 305, Japan
⁵⁴Tufts University, Medford, Massachusetts 02155, USA
⁵⁵University of Virginia, Charlottesville, Virginia 22906, USA
⁵⁶Waseda University, Tokyo 169, Japan
⁵⁷Wayne State University, Detroit, Michigan 48201, USA
⁵⁸University of Wisconsin, Madison, Wisconsin 53706, USA
⁵⁹Yale University, New Haven, Connecticut 06520, USA

We present a measurement of the top-quark mass using a sample of $t\bar{t}$ events in 5.7 fb^{-1} of integrated luminosity from $p\bar{p}$ collisions at the Fermilab Tevatron with $\sqrt{s} = 1.96 \text{ TeV}$ and collected by the CDF II Detector. We select events having no identified charged leptons, large missing transverse energy, and four, five, or six jets with at least one jet tagged as coming from a b quark. This analysis considers events from the semileptonic $t\bar{t}$ decay channel, including events that contain tau leptons. The measurement is based on a multidimensional template method. We fit the data to signal templates of varying top-quark masses and background templates, and measure a top-quark mass of $M_{\text{top}} = 172.3 \pm 2.4 \text{ (stat)} \pm 1.0 \text{ (syst)} \text{ GeV}/c^2$.

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*Deceased

[†]With visitors from ^aIstituto Nazionale di Fisica Nucleare, Sezione di Cagliari, 09042 Monserrato (Cagliari), Italy, ^bUniversity of CA Irvine, Irvine, CA 92697, USA, ^cUniversity of CA Santa Barbara, Santa Barbara, CA 93106, USA, ^dUniversity of CA Santa Cruz, Santa Cruz, CA 95064, USA, ^eCERN, CH-1211 Geneva, Switzerland, ^fCornell University, Ithaca, NY 14853, USA, ^gUniversity of Cyprus, Nicosia CY-1678, Cyprus, ^hOffice of Science, U.S. Department of Energy, Washington, DC 20585, USA, ⁱUniversity College Dublin, Dublin 4, Ireland, ^jUniversity of Fukui, Fukui City, Fukui Prefecture, Japan 910-0017, ^kUniversidad Iberoamericana, Mexico D.F., Mexico, ^lIowa State University, Ames, IA 50011, USA, ^mUniversity of Iowa, Iowa City, IA 52242, USA, ⁿKinki University, Higashi-Osaka City, Japan 577-8502, ^oKansas State University, Manhattan, KS 66506, USA, ^pUniversity of Manchester, Manchester M13 9PL, United Kingdom, ^qQueen Mary, University of London, London, E1 4NS, United Kingdom, ^rUniversity of Melbourne, Victoria 3010, Australia, ^sMuons, Inc., Batavia, IL 60510, USA, ^tNagasaki Institute of Applied Science, Nagasaki, Japan, ^uNational Research Nuclear University, Moscow, Russia, ^vUniversity of Notre

The top quark (t) is the heaviest known elementary particle, with a mass approximately 40 times heavier than the mass of the bottom quark (b), its partner in the weak isospin doublet. The top-quark mass can be used as a consistency check for standard model (SM) parameters. For example, the top quark contributes significantly to electroweak radiative corrections relating its mass (M_{top}) and that of the W boson to the mass of the predicted Higgs boson within the SM and in new physics models [1, 2]. Therefore, combined precision measurements of the W mass and M_{top} provide an important

Dame, Notre Dame, IN 46556, USA, ^wUniversidad de Oviedo, E-33007 Oviedo, Spain, ^xTexas Tech University, Lubbock, TX 79609, USA, ^yUniversidad Tecnica Federico Santa Maria, 110v Valparaíso, Chile, ^zYarmouk University, Irbid 211-63, Jordan, ^{hh}On leave from J. Stefan Institute, Ljubljana, Slovenia,

constraint on the Higgs boson mass. At the Tevatron, the top quark is predominantly produced in $t\bar{t}$ pairs. As in the SM the top quark decays almost exclusively to a W boson and a b quark, the expected signature of a $t\bar{t}$ production event is $t\bar{t} \rightarrow W^+bW^-\bar{b}$, assuming unitarity of the three-generation CKM matrix [3]. Since the W boson subsequently decays either to a quark-antiquark pair or to a lepton-neutrino pair, the final state of $t\bar{t}$ production can be classified by the number of charged leptons produced. In this letter, we focus on events with large missing transverse energy (\cancel{E}_T) [4] as expected for undetected energetic neutrinos, accompanied by jets. We explicitly veto events with identified high p_T electrons or muons (lepton+jets events) as well as multijet events where both W bosons decay hadronically (all-hadronic events). This ensures that our result is statistically independent from other CDF top-quark mass measurements [5–8] and allows for a future combination with them. A previous measurement of M_{top} in this final state used an integrated luminosity of 311 pb^{-1} [9] and yielded $M_{\text{top}} = 172.3 \pm 15.3 \text{ (stat)} \pm 14.4 \text{ (syst)} \text{ GeV}/c^2$. Although no identified leptons are explicitly required, our measurement is sensitive to all W leptonic decays. This includes decays to τ leptons, which constitute approximately 40% of the signal sample in our final selection. The frequency of observed τ lepton final states is predicted to be enhanced by new physics models such as a charged Higgs decay [10], therefore, a significantly different measurement of M_{top} in this decay channel could indicate contributions from non-SM [11] physics processes.

We use data corresponding to an integrated luminosity of 5.7 fb^{-1} of $p\bar{p}$ collisions at the Fermilab Tevatron, and collected by the CDF II Detector [12]. The sample of events used in this measurement is a subset of events that initially passed a trigger requirement, which accepted events with at least four calorimetric clusters [13] of $E_T > 15 \text{ GeV}$ and a scalar sum E_T of these clusters greater than 175 GeV [14]. After the trigger selection, event observables of physical interest are computed. Jets are reconstructed with the JETCLU [15] algorithm using a cone radius of $\Delta R = \sqrt{\Delta\eta^2 + \Delta\phi^2} = 0.4$ [16]. Jets are corrected [17] for nonuniformities of the calorimeter response as a function of η , energy contributed by multiple $p\bar{p}$ interactions in the event, and calorimeter nonlinear response. To determine if a jet originated from a b quark, a secondary vertex algorithm [18] is applied. This algorithm identifies jets that are likely to come from b quark hadronization through the presence of a displaced vertex within the jets (b -tag). We require at least one jet to be identified as a b jet (b -tagged). We divide the sample of candidate events into two, separating events with one b -tagged jet (1-tag) from events with two or more b -tagged jets (2-tag). Events are required to have at least four and at most six jets with transverse energy $E_T > 15 \text{ GeV}$ and $|\eta| < 2.0$. To avoid overlap with other CDF top-quark mass measurements, we reject events with reconstructed electrons or muons with $p_T > 20 \text{ GeV}/c$ and $|\eta| < 1.0$ (lepton + jets final state),

and events with \cancel{E}_T significance below $3 \text{ GeV}^{1/2}$ (all-hadronic final state), where the \cancel{E}_T significance is defined as $\cancel{E}_T^{\text{sig}} = \cancel{E}_T / \sqrt{\sum_{\text{jets}} E_T}$. For further rejection of multijet backgrounds from QCD processes, we require $\Delta\phi_{\text{min}}(\cancel{E}_T, \text{jet}) > 0.4$, where $\Delta\phi_{\text{min}}(\cancel{E}_T, \text{jet})$ is the smallest separation in the angle ϕ between jets and \cancel{E}_T [14].

Background events with b -tags arise from QCD multijet and electroweak production of W bosons associated with heavy flavor jets. In order to improve the signal-to-background ratio (S/B) in this analysis, an artificial neural network is trained to identify the kinematic and topological characteristics of SM $t\bar{t}$ events using eight input variables [14]. We apply the neural network to all events passing the above selections, and make a cut on neural network score that retains 81% of the $t\bar{t}$ signal events while rejecting 91% of background events. The selection criteria listed above define the ‘signal region’. We follow [14] and estimate the background rate using a data-driven method. The method uses events with exactly three jets and employs a per-jet parameterization of the b -tagging probability. Due to the presence of $t\bar{t}$ events in samples with higher jet multiplicity, we extrapolate the b -tagging probability of the three-jet event sample to higher jet multiplicity events by iteratively removing the $t\bar{t}$ content from the sample [14]. We estimate the background for the 1-tag and 2-tag samples separately. A b -tagging correction factor [19] is applied to take into account the fact that most of the heavy flavor jets are produced in pairs. With this procedure we obtain the estimated number of background events in the signal region shown in Table I. We also show the estimated number of $t\bar{t}$ signal events, assuming a $t\bar{t}$ production cross-section of 7.5 pb at $M_{\text{top}} = 172.5 \text{ GeV}/c^2$ [20], together with the number of observed events in the data.

Monte Carlo (MC)-simulated $t\bar{t}$ samples are generated by PYTHIA [21] with 76 different M_{top} values ranging from $150 \text{ GeV}/c^2$ to $240 \text{ GeV}/c^2$. For each of them we reconstruct the events with different values of Δ_{JES} , the difference between the jet energy scale (JES) in the MC simulation and the data. A total of 27 different Δ_{JES} values, ranging from $-3.0 \sigma_c$ to $+3.0 \sigma_c$, where σ_c is the uncertainty on the JES [22], are used to reduce the systematic effects due to the jet energy uncertainty, as described later.

For each MC-simulated sample in this analysis, we reconstruct three variables using the leading four or five jets to form the templates. The first variable, m_{jj} , defined as the invariant mass of the two jets from the hadronically decaying W boson, serves as an *in situ* constraint of the JES through the likelihood fit described later. We calculate m_{jj} from the two non- b -tagged jets whose invariant mass produces the closest value to the world average W boson mass of $80.40 \text{ GeV}/c^2$ [23]. We also extract the reconstructed top-quark mass ($m_t^{\text{had, reco}}$) from the invariant mass of the three jets whose momentum sum yields the largest p_T since the invariant mass constructed this

TABLE I: Number of expected signal and background events and observed data events with integrated luminosity of 5.7 fb^{-1} in the signal region. The $t\bar{t}$ signal contribution is estimated with a cross-section of 7.5 pb . All selection requirements are applied, and events are separated into 1-tag and 2-tag categories.

Events	1-tag	2-tag
Expected $t\bar{t}$ Signal	644.3 ± 118.7	262.9 ± 50.3
Expected Background	410.6 ± 31.7	43.8 ± 11.0
Total Expectation	1054.9 ± 122.9	306.7 ± 51.5
Observed Data	1147	285

way has a large correlation with the hadronically decaying top-quark mass. To enhance the probability that the jets used to compute $m_t^{\text{had, reco}}$ come from the hadronically decaying top quark, we add two constraints to the calculation of $m_t^{\text{had, reco}}$: first, $m_t^{\text{had, reco}}$ must contain the two jets that contribute to m_{jj} , and second, in 2-tag events one jet of $m_t^{\text{had, reco}}$ must be b -tagged. A third variable, $m_t^{\text{had, reco}(2)}$, is defined as the invariant mass of three jets, two of which are the pair which defines m_{jj} . The third jet of $m_t^{\text{had, reco}(2)}$ is required to be the most energetic jet of those not forming $m_t^{\text{had, reco}}$, and is required to be b -tagged in a 2-tag event. The variable $m_t^{\text{had, reco}(2)}$ provides complementary information on the top quark mass, and is particularly important in events where the three jets used to compute $m_t^{\text{had, reco}}$ were not the actual decay products of the hadronically-decaying top quark.

The template method used in the extraction of M_{top} requires that a probability density function (p.d.f.) be built for each template. For each MC signal and background sample, we estimate the p.d.f.s using the kernel density estimation (KDE)[24, 25] that employs a non-parametric method to construct p.d.f.s. For each sample, we build a three dimensional p.d.f. from the reconstructed observables (m_{jj} , $m_t^{\text{had, reco}}$, and $m_t^{\text{had, reco}(2)}$), taking their correlations into account. To measure M_{top} , we fit the signal and background p.d.f.s to the distributions of the observables in the data using an unbinned maximum likelihood fit [26] where we minimize the negative logarithm of the likelihood using MINUIT [27]. The likelihood fits M_{top} and Δ_{JES} simultaneously, and is built separately for each subsample, 1-tag and 2-tag events, in order to improve the usage of statistical information. References [25, 28] provide detailed information about this technique.

The mass fitting procedure is tested with pseudoexperiments for a set of MC $t\bar{t}$ samples with 14 different M_{top} values ranging from $159 \text{ GeV}/c^2$ to $185 \text{ GeV}/c^2$. For each pseudoexperiment we select the number of background events from a Poisson distribution with a mean equal to the expected total number of background events in the sample and the number of signal events from a Poisson distribution with a mean equal to the expected number of signal events normalized to a $t\bar{t}$ production cross

section of 7.5 pb [20]. The distributions of the average mass residual (the deviation from the input top-quark mass) and the width of the pull (the ratio of the residual to the uncertainty) for simulated experiments are corrected to be zero and unity, respectively. The correction is $M_t^{\text{corr}} = 1.24 \times M_t^{\text{meas}} - 40.6 \text{ GeV}/c^2$, where M_t^{meas} is the raw value from the likelihood fit and M_t^{corr} is the corrected value of the measurement. The measured uncertainty is correspondingly increased by 20% to correct the width of the pull distribution.

We examine various sources of systematic uncertainties that could affect the measurement by comparing the results of pseudoexperiments in which we vary relevant parameters within their uncertainties. One of the dominant sources of systematic uncertainty is the residual JES [6, 22]. We vary the JES components within their uncertainties in the generated signal MC events and interpret the shifts in the returned top-quark mass as uncertainties. The b jet energy scale systematic uncertainty that arises from the modeling of b fragmentation, b hadron branching fractions, and calorimeter response captures the additional uncertainty not taken into account in the light quark jet energy scale [6]. The uncertainty arising from the choice of MC generator is estimated by comparing results from MC simulated samples generated with PYTHIA and HERWIG [29]. We estimate the systematic uncertainty due to imperfect modeling of initial-state gluon radiation (ISR) and final-state gluon radiation (FSR) by varying the amounts of ISR and FSR in simulated events [30]. We estimate the systematic uncertainty due to parton distribution functions (PDF's) of the proton by varying the independent eigenvectors of the CTEQ6M [31] PDF's, varying Λ_{QCD} (228 MeV vs. 300 MeV), and comparing CTEQ5M [32] with MRST72 [33] PDF's. To estimate the systematic uncertainty associated with uncertainties in the top-quark production mechanism, we vary the fraction of the top quarks produced by gluon-gluon annihilation from the default 6% to 20%, corresponding to a one standard deviation upper bound on the gluon fusion fraction [34]. We also evaluate the uncertainty due to background modeling effects by re-weighting the background shape up and down and comparing the resulting measurements. We apply an additional uncertainty to account for the effect of the trigger simulation in the signal MC samples, in a similar way to the background shape systematic uncertainty estimation. We also estimate an uncertainty due to the effect of multiple hadron interactions, which takes into account the increasing instantaneous luminosity in this dataset. The color reconnection (CR) systematic uncertainty [35] is evaluated using MC samples generated with and without CR effects adopting different tunes [36] of PYTHIA. Table II summarizes the individual systematic uncertainties considered, giving a total systematic uncertainty of $1.0 \text{ GeV}/c^2$ for the measurement of M_{top} .

By applying a likelihood fit to the data using the three observables described above and the corrections obtained from the simulated experiments, the top-quark mass is

TABLE II: Systematic uncertainties on M_{top} .

Systematic Sources	Uncertainty (GeV/c^2)
Residual JES	0.5
b -JES	0.3
MC Generator	0.7
Radiation	0.2
PDFs	0.2
gg fraction	< 0.1
Background	0.1
Trigger Modeling	0.1
Multiple Hadron Interaction	0.2
Color Reconnection	0.2
Total Systematic Uncertainty	1.0

measured to be

$$\begin{aligned}
M_{\text{top}} &= 172.3 \pm 2.4(\text{stat}) \pm 1.0(\text{syst}) \text{ GeV}/c^2 \\
&= 172.3 \pm 2.6 \text{ GeV}/c^2.
\end{aligned} \tag{1}$$

Figure 1 shows the distribution of the observables used for the M_{top} measurement overlaid with their probability density functions from $t\bar{t}$ signal events with $M_{\text{top}} = 172.5 \text{ GeV}/c^2$ and the estimated background.

In conclusion, we have performed a measurement of the top-quark mass in events with large \cancel{E}_T and jets, corre-

sponding to an integrated luminosity of 5.7 fb^{-1} . The data sample has been chosen in such a way as to exclude events used in other CDF top-quark mass measurements. The result, $M_{\text{top}} = 172.3 \pm 2.6 \text{ GeV}/c^2$, is approximately a factor of 6 improvement from the previous measurement in this channel [9], and is in agreement with other measurements which contribute to the world average of $M_{\text{top}} = 173.2 \pm 0.9 \text{ GeV}/c^2$ [37].

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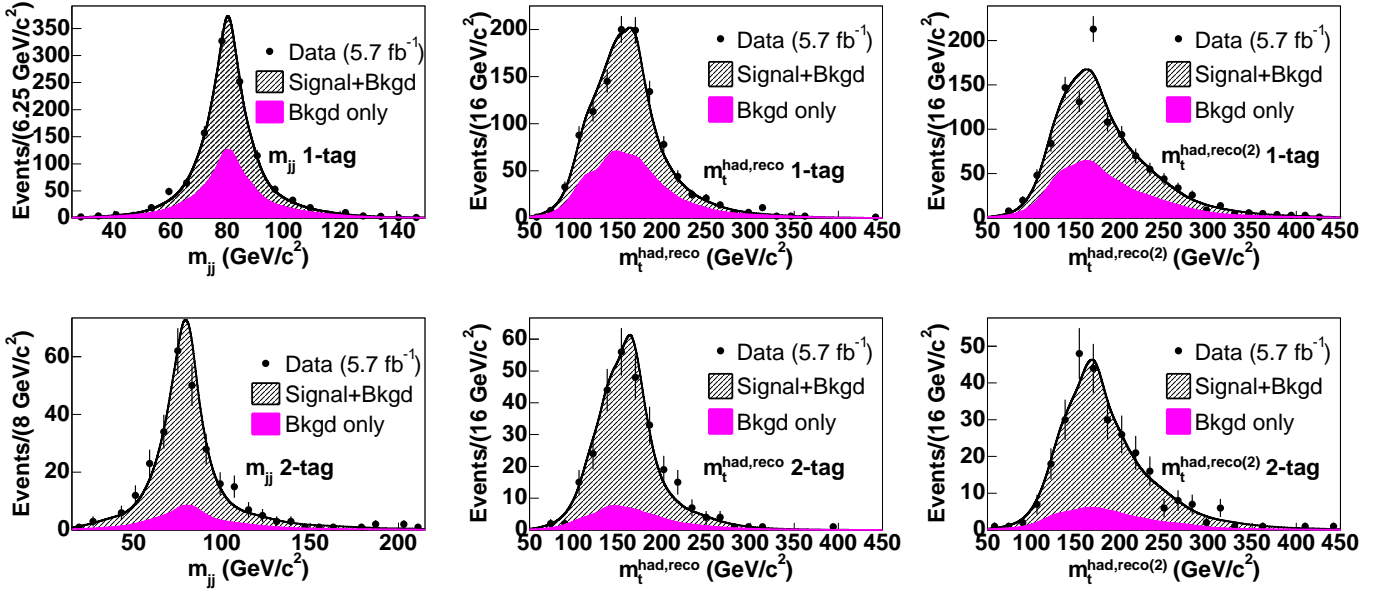


FIG. 1: Distribution of kinematic variables m_{jj} , $m_t^{\text{had,reco}}$, and $m_t^{\text{had,reco}(2)}$ from data (points), overlaid with their corresponding one dimensional p.d.f.s from signal MC sample ($M_{\text{top}} = 172.5 \text{ GeV}/c^2$, hashed area) plus the estimated background (filled area). Both 1-tag (top) and 2-tag (bottom) events are displayed.

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